

DIRECT FORMULA FOR \bar{E} in H.O

$\bar{E} = \frac{1}{2} kT$ for each term (K & D P.E)

c.g. In 3-D, 3 K.E & 3 P.E terms

$\bar{E} = kT + kT + kT = 3kT$

$\bar{E} = 3kT$

but if particles are non interacting

$\Rightarrow \bar{E} = \frac{1}{2} kT + \frac{1}{2} kT + \frac{1}{2} kT$

$\Rightarrow \bar{E} = \frac{3}{2} kT$

For H.O.

$Z = \sum_{n=0}^{\infty} e^{-\beta (n + \frac{1}{2}) h\nu}$

$Z = e^{-h\nu\beta/2} + e^{-3h\nu\beta/2} + \dots$

GP series

$a = e^{-h\nu\beta/2} \quad r = e^{-\beta h\nu}$

$\therefore Z = \frac{e^{-\beta h\nu/2}}{1 - e^{-\beta h\nu}}$

For 3D HO, $Z_N = \left(\frac{1}{1 - e^{-\beta h\nu}}\right)^3 N$

Average value of $n = \frac{1}{2}$ (classical)

Partition fn for particles in magnetic field

$Z = e^{-\beta E} = e^{-\beta(\vec{\mu} \cdot \vec{H})} = e^{-\beta(\vec{\mu} \cdot \vec{B})}$

up $\rightarrow \mu B$, down $\rightarrow -\mu B$

$Z = e^{-\mu\beta B} + e^{\mu\beta B}$

EX JUNE 2012

EXAMPLE \rightarrow If $F = NkT \ln\left(\frac{2\omega V}{N} (kT)^{5/2}\right)$

Find U.

Soln. As $F = -NkT \ln Z$

$\Rightarrow Z = \frac{2\omega V (kT)^{5/2}}{N}$

$U = -\frac{\partial}{\partial \beta} (\ln Z) = -\frac{\partial}{\partial \beta} \left(\ln\left(\frac{2\omega V}{N} (kT)^{5/2}\right) \right)$

$= \frac{5N}{2\beta} = \frac{5}{2} NkT$

$U = \frac{5}{2} NkT$

Ensemble \rightarrow system of collection of similar systems having some const parameters

Microcanonical

Walls are rigid impermeable well insulated

N, V, E

Microcanonical

Canonical

Walls are rigid conducting impenetrable

V, N, T

Canonical

Grand canonical

Walls are impenetrable rigid conducting

μ, V, T

Grand canonical

These quantities remain same in a particle ensemble i.e. in microcanonical, canonical & Grand canonical resp. $\frac{1}{kT} \leftarrow$ Energy fluctuation $\rightarrow \frac{1}{N}$

Maxwell Boltzmann distⁿ law of velocities

No. of particles in velocity range v to $v+dv$ is

$n(v_x) dv_x = N \left(\frac{m}{2\pi kT}\right)^{1/2} e^{-\frac{mv_x^2}{2kT}} dv_x$

$P = \frac{n}{N} = \left(\frac{m}{2\pi kT}\right)^{1/2} e^{-\frac{mv_x^2}{2kT}}$

$P_{max} = \left(\frac{m}{2\pi kT}\right)^{1/2} \left(\frac{kT}{m}\right)^{1/2}$

$\bar{v}_x = \int_{-\infty}^{\infty} v_x P dv_x = 0$

$\bar{v}_x = \bar{v}_y = \bar{v}_z = 0$

$\langle v_x v_y \rangle = \langle v_y v_x \rangle = \langle v_z v_x \rangle = 0$

In 1-D $\langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle = kT/m$

In 2-D $\langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle = 2kT/m$

In 3-D $\langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle = 3kT/m$

Statistical Physics

EW Minium



Statistical Physics:

This book delves into Statistical Physics. Statistical Physics is an essential topic that needs to be grasped by everyone, ranging from students and scholars to the general public. This book will furnish comprehensive and in-depth insights into Statistical Physics, encompassing both the fundamentals and more intricate discussions.

1. This book is structured into several chapters, namely:

- Chapter 1: Introduction to Statistical Physics
- Chapter 2: Essential Elements of Statistical Physics
- Chapter 3: Statistical Physics in Everyday Life
- Chapter 4: Statistical Physics in Specific Contexts
- Chapter 5: Conclusion

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4. In chapter 3, the author will examine the practical applications of Statistical Physics in daily life. The third chapter will showcase real-world examples of how Statistical Physics can be effectively utilized in everyday scenarios.

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The book is crafted in an easy-to-understand language and is complemented by engaging illustrations. This book is highly recommended for anyone seeking to gain a comprehensive understanding of Statistical Physics.

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web it stands at 17 metres tall there aren t any real honeybees around it but many humans find themselves sucked into the hive at kew it is a striking insta

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